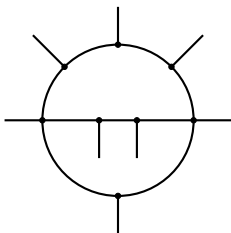


Motivic geometry of two-loop Feynman integrals



Andrew Harder
Lehigh University

Joint with Chuck Doran and Pierre Vanhove

Feynman graphs

Suppose Γ is a graph, $D \in 2\mathbb{Z}_{\geq 0}$ is space-time dimension.

1. $E_\Gamma, V_\Gamma, H_\Gamma$, edges, vertices and half (external) edges
2. To each half-edge, we associate a momentum $p_h \in \mathbb{R}^{1,D-1}$, satisfying;

$$\sum_{h \in H_\Gamma} p_h = 0 \quad (\text{conservation of momentum})$$

3. To each internal edge we associate a mass $m_e \in \mathbb{R}_{\geq 0}$.

Simplifications:

- View masses and momenta as complex numbers.
- No 1-valent vertices.
- For each vertex attach a single half-edge (write p_v instead of p_h).

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Consider the ring $\mathbb{C}[x_e \mid e \in E_\Gamma]$.

U_Γ polynomial

Suppose T is a spanning tree of Γ ,

$$x^T = \prod_{e \notin T} x_e, \quad U_\Gamma = \sum_{\text{spanning trees}} x^T.$$

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Suppose $T_1 \cup T_2$ is a spanning 2-tree of Γ .

$$x^{T_1 \cup T_2} = \prod_{e \notin T_1 \cup T_2} x_e, \quad V_\Gamma = \sum_{\text{spanning 2-trees}} s_{T_1 \cup T_2} x^{T_1 \cup T_2}$$

Here $s_{T_1 \cup T_2} = (\sum_{v \in T_1} p_v)^2 = (\sum_{v \in T_2} p_v)^2$.

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$$\mathbf{F}_\Gamma = \mathbf{U}_\Gamma \left(\sum_e m_e^2 x_e \right) + \mathbf{V}_\Gamma.$$

$$\deg \mathbf{U}_\Gamma = \ell(\Gamma), \quad \deg \mathbf{F}_\Gamma = \ell(\Gamma) + 1$$

Example:



$$\mathbf{U}_\Gamma = x_0 x_1 + x_1 x_2 + x_0 x_2$$

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Remark: In nearly *all* examples, $V(\mathbf{F}_\Gamma)$ is *singular*.

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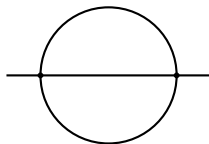
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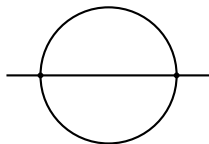
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$$I_{\Gamma}(m, s) = \int_{\sigma} \frac{\mathbf{U}_{\Gamma}^{\epsilon_{\Gamma} - D(\ell+1)/2}}{\mathbf{F}_{\Gamma}^{\epsilon_{\Gamma} - D\ell/2}} \Omega_0$$

$$\sigma = \{[x_1 : \cdots : x_{\epsilon_{\Gamma}}] \in \mathbb{P}^{\epsilon_{\Gamma}-1} \mid x_i \in \mathbb{R}_{\geq 0}\}$$

$$\Omega_0 = \sum_{i=1}^{\epsilon_{\Gamma}} (-1)^i x_i \left(dx_1 \wedge \cdots \wedge \widehat{dx}_i \wedge \cdots \wedge dx_{\epsilon_{\Gamma}} \right)$$

Periods and mixed Hodge structures

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where $Z_{\Gamma,D} = V(\mathbf{F}_\Gamma)$, $V(\mathbf{F}_\Gamma \mathbf{U}_\Gamma)$, or $V(\mathbf{U}_\Gamma)$, and $B_\Gamma = V\left(\prod_{i=0}^{\text{er}} x_i\right)$

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$$[\sigma] \notin \mathbf{H}_{\text{er}-1}(\mathbb{P}^{\text{er}-1} - Z_{\Gamma,D}, B_\Gamma) \text{ since } \sigma \cap Z_{\Gamma,D} \neq \emptyset$$

Theorem: (Bloch–Esnault–Kreimer, Brown) After appropriate (toric) blow up $b : \mathbb{P}_\Gamma \rightarrow \mathbb{P}^{\text{er}-1}$ and modification of σ , the Feynman integral is a well-defined relative period.

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Extensions of mixed Hodge structure

The mixed Hodge structure on $H^{\text{er}-1}(\mathbb{P}_\Gamma - Z_\Gamma, B_\Gamma)$ is (up to mixed Tate factors) an iterated extension of mixed Hodge structures of the shape

$$H^*(\mathbb{P}_{\Gamma-\Gamma'} - Z_{\Gamma-\Gamma',D}; \mathbb{Q}), \quad H^{\text{er} // \Gamma' - 1}(\mathbb{P}_{\Gamma // \Gamma'} - Z_{\Gamma // \Gamma'}; \mathbb{Q})$$

Goal. We want to study the mixed Hodge structures on the cohomology groups of $Z_{\Gamma,D}$.

(I.e., we do not study Feynman integrals individually, rather we study a space of functions in which they live).

From here on, will study the case where $Z_{\Gamma,D} = V(\mathbf{F}_\Gamma)$.

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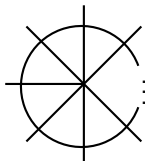
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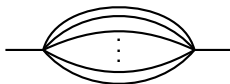
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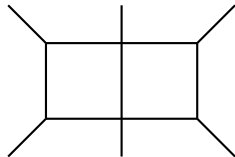
Wheel



Sunset family



Double box

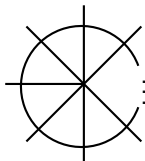


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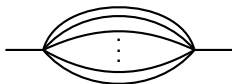
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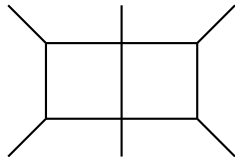
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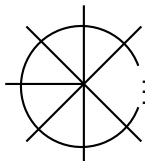


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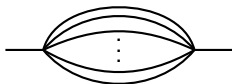
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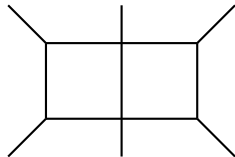
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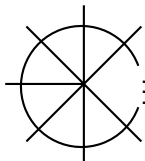


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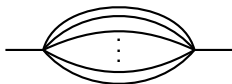
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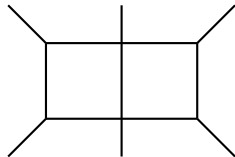
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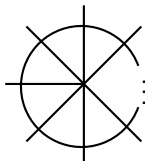


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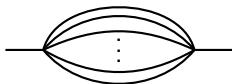
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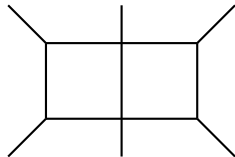
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“Proposition”. Suppose $\pi : Z \rightarrow \mathbb{P}^n$ is a quadric fibration of dimension n with generic corank r and let \mathcal{D} be a divisor along which the corank of quadrics increases.

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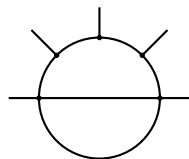
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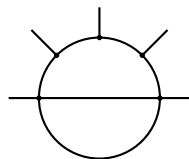
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Possible cases:

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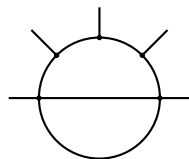
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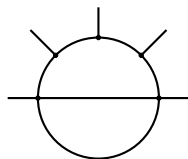
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Example: $\Gamma_{a,1,1}$ graphs



||

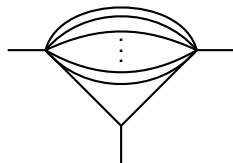
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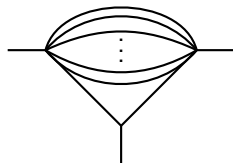
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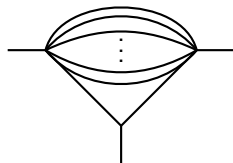
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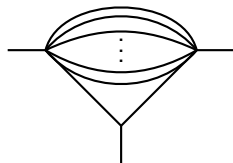
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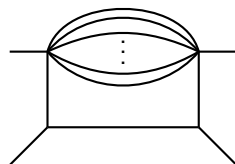
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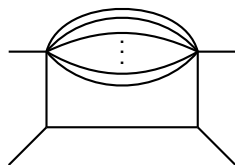
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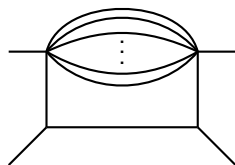
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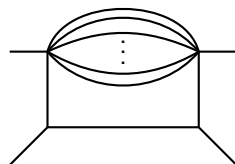
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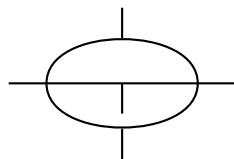
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The discriminant locus is a nodal quartic surface along with a hyperplane.

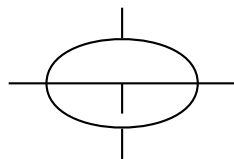
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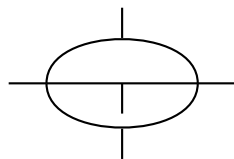
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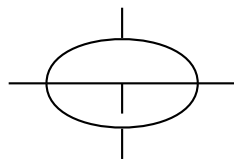
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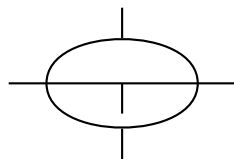
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Feynman cubics ($\ell = 2$)

$$\Gamma_{a,b,c} = \text{Diagram} \quad \begin{array}{l} \deg \mathbf{F}_{\Gamma_{a,b,c}} = 3 \\ \deg \mathbf{U}_{\Gamma_{a,b,c}} = 2 \end{array}$$

The MHS attached to $\mathbf{U}_{\Gamma_{a,b,c}}$ is mixed Tate (“easy”) so we will try to understand the case where $Z_{\Gamma_{a,b,c},D} = V(\mathbf{F}_{\Gamma_{a,b,c}})$.

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The main result (planar two-loop graphs)

Theorem (Doran-H.-Vanhove). Suppose that $b = 0, 1$ (that is, $\Gamma_{a,b,c}$ is a planar graph). Then if $H^*(V(\mathbf{F}_{\Gamma_{a,b,c}}); \mathbb{Q})$ is contained in

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If a or $c \leq 2$ then we can replace “hyperelliptic” with “elliptic”. If dimension is large enough relative to a and c then cohomology is mixed Tate.

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Proposition. If a cubic hypersurface X contains a linear subspace L of codimension 1, then $H^*(X; \mathbb{Q}) \in \mathbf{MHS}_{\mathbb{Q}}^{\text{hyp}}$.

Proof. Take the pencil of hyperplane sections in X containing L . This induces a quadric fibration over \mathbb{P}^1 . The cohomology of a quadric fibration over \mathbb{P}^1 is either (a) mixed Tate (if relative dimension is even) or hyperelliptic (if relative dimension is odd). □

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Proof sketch (continued)

There is a birational map: $\phi : V(\mathbf{F}_{\Gamma_{a,1,c}}) \dashrightarrow X_{a,1,c}$ and open subsets

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Motivation coming from Picard–Fuchs operators

Lairez–Vanhove: For many 2-loop Feynman graphs, compute the Picard–Fuchs operators of the family of hypersurfaces

$$\mathbf{F}_\Gamma(t) = \mathbf{U}_\Gamma \left(\sum_{e \in E_\Gamma} m_e^2 x_e \right) + t \mathbf{V}_\Gamma, \quad \omega_\Gamma = \frac{\mathbf{U}_\Gamma^{\text{er}-D(\ell+1)/2}}{\mathbf{F}_\Gamma^{\text{er}-D\ell/2}} \Omega$$

I.e., find a differential operator $\mathcal{L}_\Gamma \in \mathbb{C}[t, \partial_t]$ so that

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Remark: $l_\Gamma(t)$ satisfies an inhomogeneous ODE

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Observation: Frequently, \mathcal{L}_Γ factors into components which are

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$\text{Sol}(\mathcal{L}_\Gamma)$ is a subquotient of $\mathcal{H}_\Gamma = \mathbb{C}$ -VMHS attached to $V(\mathbf{F}_\Gamma(t))$.

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