Tameness in Hodge Theory and Physics

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Based on:

2302.04275 Part II

2210.10057 Part I

2112.08383

2112.06995

with Michael Douglas, Lorenz Schlechter

TG

with Ben Bakker, Christian Schnell, Jacob Tsimerman

Motivation

 Much recent activity in mapping out the tame parts of mathematics (algebraic geometry, arithmetic geometry, number theory,...)

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 - Tameness of the period map, mixed period map, period integrals

 [Bakker, Klingler Tsimerman '18]...
 - Ax-Schanuel conjecture for Hodge structures [Bakker, Tsimerman '17]... several subsequent generalizations, e.g. to mixed Hodge structures
 - · Finiteness of self-dual integral classes [Bakker, TG, Schnell, Tsimerman '21]
 - Geometric André-Grothendieck Period Conjecture [Bakker, Tsimerman '22]

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- consider 12-dimensional F-theory: solving Einstein's equations and other equations of motion

 - · 12d manifold: $\mathbb{S} \times Y$ Calabi-Yau fourfold · 4-form: $G_4 \in H^4(Y,\mathbb{Z})$ $\int_Y G_4 \wedge G_4 = \ell$ * $G_4 = G_4$ (in cohom.)
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- Conjecture [Douglas '03] [Acharya, Douglas '06]:
 - Number of distinct solutions of string theory with bounds on vacuum energy, KK scale, compactification volume are finite
 - \rightarrow in the above setting: finitely many choices for G_4 .

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- → C be Weil operator: $Cv = i^{p-q}v$ $v \in H^{p,q}$ (e.g. Hodge star C = *) $Q \text{ be polarization: } Q(\bar{v}, Cv) > 0$ (e.g. $Q(v, w) = \int_{V} v \wedge w$)

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Theorem [Bakker,TG,Schnell,Tsimerman]: For integer $\ell > 0$, the locus of integral self-dual classes

$$S_{\ell} = \left\{ (x, v) \in E : v \in H_{\mathbb{Z}, x} \text{ and } C_x v = v \text{ and } Q(v, v) = \ell \right\}$$

is a set definable in the o-minimal structure $\mathbb{R}_{\mathrm{an,exp}}$

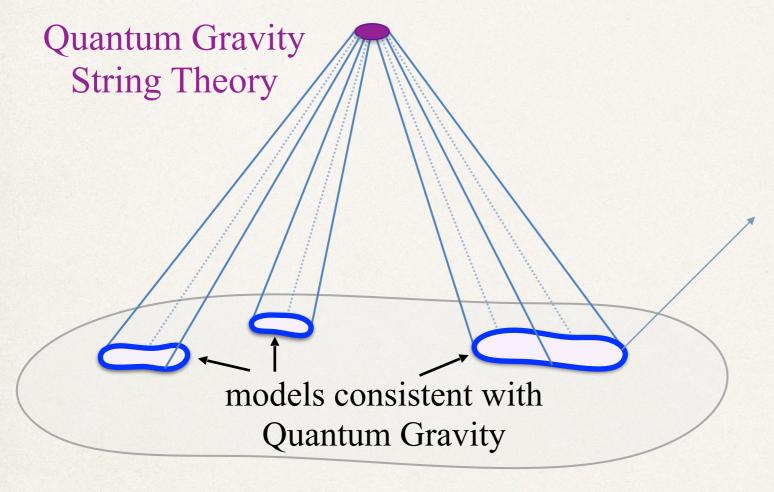
 \rightarrow \mathcal{S}_{ℓ} has finitely many connected components

Tameness in Physics

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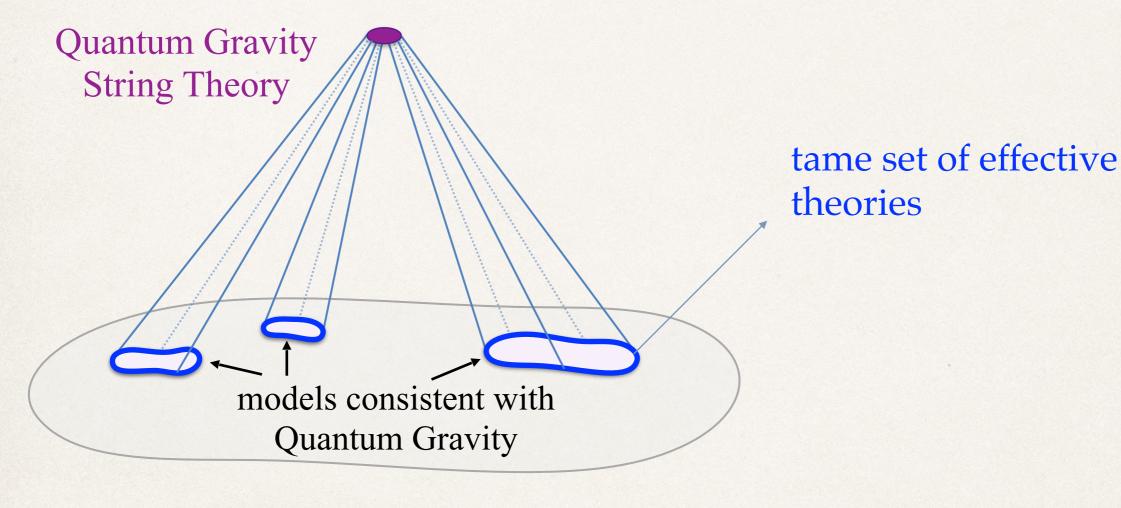
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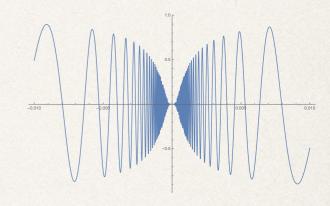
- Tameness in Quantum Field Theories (QFTs) [Douglas,TG,Schlechter '22+'23]
 - → use of many of the recent results on tameness in Hodge theory
 - → several new conjectures

Tameness and o-minimal structures

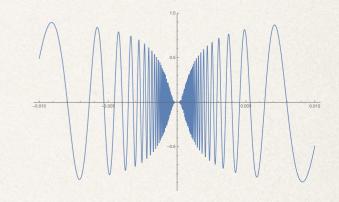
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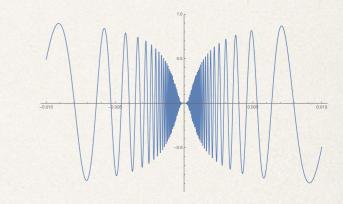


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- Grothendieck's dream:
 develop mathematical framework for geometry
 - → tame topology [Esquisse d'un programme]

Tameness - Definition

- structure S: collect subsets of \mathbb{R}^n , n = 1, 2, ...
 - Closed under finite unions ∨, intersections ∧, complements ¬, products
 - Closed under projections (existential quantifier ∃)
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sets in o-minimal structure S: tame sets

functions with graph being a tame set: tame functions

→ tame manifold, tame bundles... a tame geometry

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- Logic perspective: $\mathbb{R}_{\mathcal{F}} = \langle \mathbb{R}; +, \cdot, -, >, \mathcal{F} \rangle$ $\mathcal{F} = \{f_1, f_2, ...\}$

all formulas using these symbols and \land , \lor , \neg , \exists , \forall

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 - Pfaffian extension: $\mathcal{P}(\tilde{\mathbb{R}})$ includes solutions to $\partial_{x_i} f = F_i(x, f(x))$ $F_i \text{ functions in o-minimal structure } \tilde{\mathbb{R}} \text{ [Speissegger '97]}$

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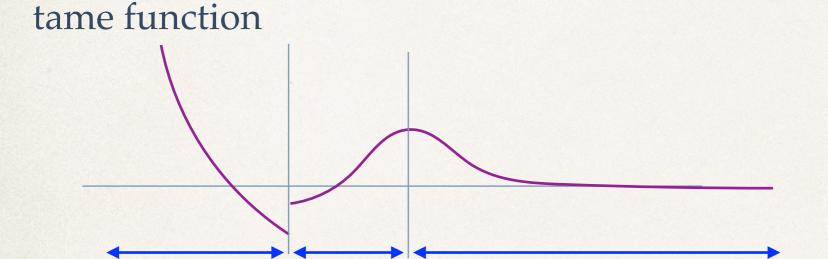
structure including $\Gamma(x)|_{(0,\infty)}$ and $\zeta(x)|_{(1,\infty)}$ [Rolin, Servi, Speissegger '22]

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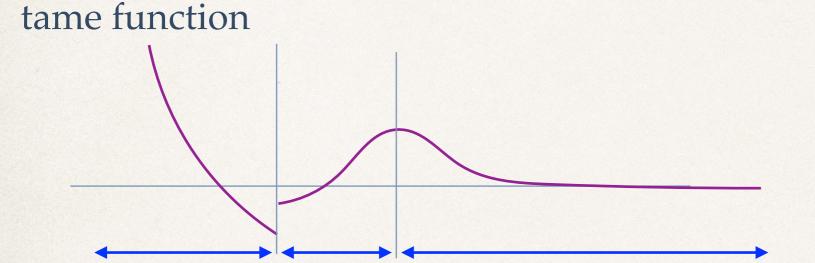
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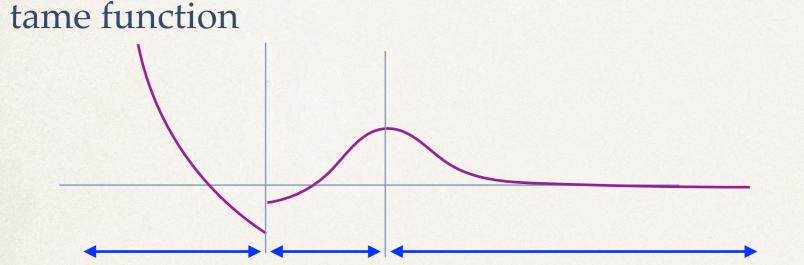


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- Periodic functions f(x+n) = f(x) are never tame (when not constant)

$$\sin(x), x \in \mathbb{R}$$



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Proof: uses crucially nilpotent orbit theorem.

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Period map is definable in $\mathbb{R}_{\mathrm{an,exp}}$.

Theorem [Cattani, Deligne, Kaplan '95]: For integer $\ell > 0$, locus of integral Hodge classes

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- original proof uses Hodge theory: nilpotent orbit theorem [Schmid]
 Sl(2) orbit theorem [Cattani, Kaplan, Schmid]

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 - → relies on holomorphicity: often absent in physical situations

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Note: - locus is, in general, only real → leave complex geometry!

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 - ⇒ works well for one-parameter limits [Schnell] [TG] '20, but too involved for multi-parameter limits

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- Proof: tameness of Weil operator period map Φ
 - tameness of maps between bundles $\Phi_E:E o \Gammaackslash (G/K imes H_{\mathbb C})$
 - lattice reduction: [e.g. Kneser] group Γ acts on set $\{v \in H_{\mathbb{Z}}: Q(v,v)=\ell\}$ with finitely many orbits.
 - tameness of self-dual locus in a single orbit

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Question 1: What are the cycles associated to self-dual classes? (like in Hodge conjecture)

→ relevant in physics 'holography' [Lüst, Vafa, Wiesner, Xu]

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Consider compact CY fourfold. Fix
$$\ell$$
. For $\dim \mathcal{M} > O(1)\ell$:
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Note: General evidence for loci realized in \mathcal{M} where period map is sl(2)-orbit.

[Graña,TG,van de Heisteeg,Herraez,Plauschinn '22]

nilpotent orbit: [TG,Monnee] in progress

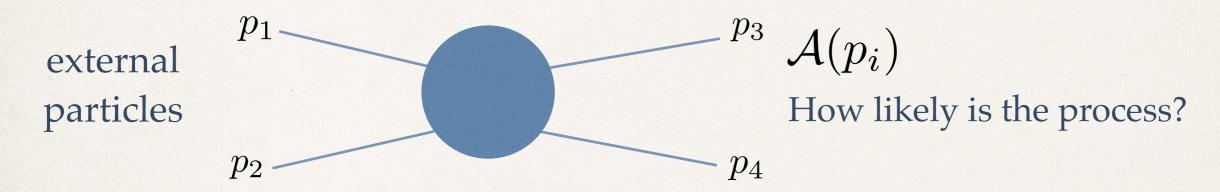
- Condition $Q(v,v)=\ell$ (tadpole condition) is central to tameness
 - → in physics this is arising from consistent coupling with gravity!
- ► A new physics conjecture 'Tadpole conjecture': see also [Hebecker's talk] [Bena,Blaback,Graña,Lüst]...[Becker,Walcher,Wrase]

Consider compact CY fourfold. Fix
$$\ell$$
. For $\dim \mathcal{M} > O(1)\ell$:
$$\min \left[\dim Comp(\mathcal{H}_{\ell})\right] > 0$$

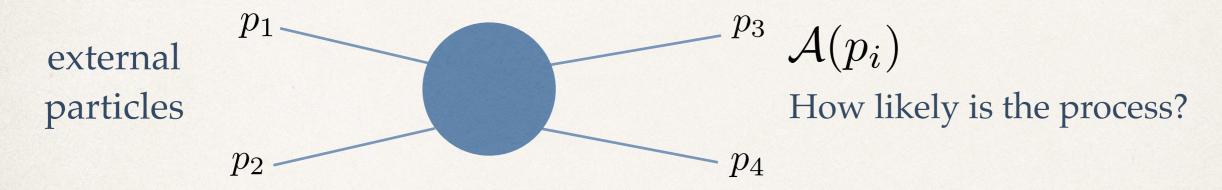
Question 2: Can one prove such a conjecture for \mathcal{H}_{ℓ} or \mathcal{S}_{ℓ} ?

Tameness in perturbative Quantum Field Theories

Scattering amplitudes

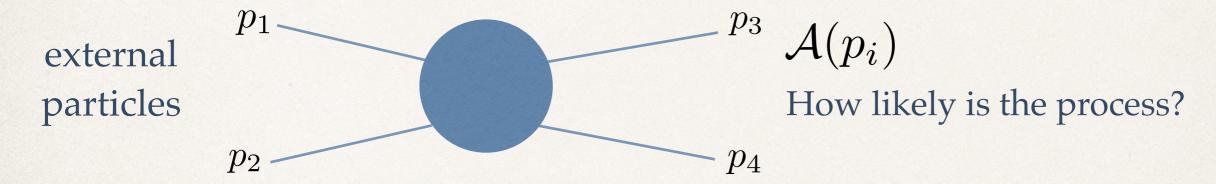


Scattering amplitudes



Physics: defined using path integrals - "sum over all possible processes"

Scattering amplitudes



- Physics: defined using path integrals "sum over all possible processes"
- Perturbative expansion: small coupling expansion $\lambda \ll 1$

→ summing till fixed loop number: finite number of Feynman integrals

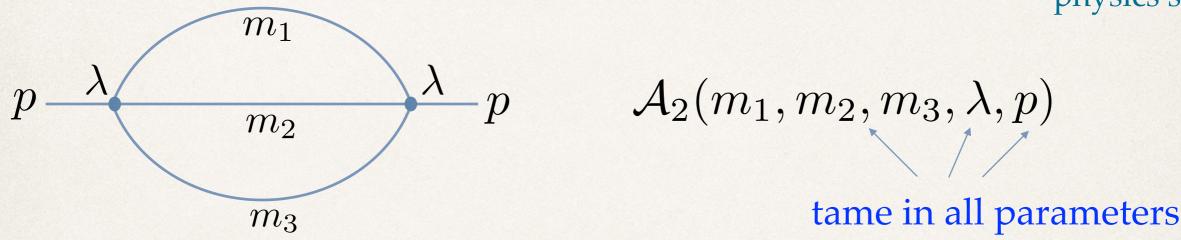
Theorem*: For any renormalizable QFT with finitely many particles and interactions all finite-loop amplitudes are tame functions of the masses, external momenta, and coupling constants.
 [Douglas,TG,Schlechter '22]

*physics style

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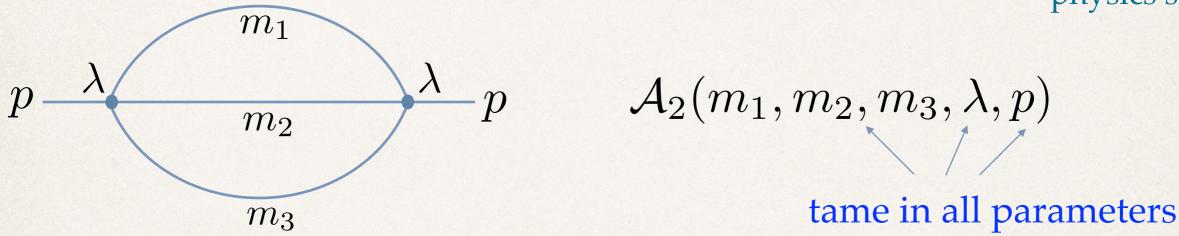
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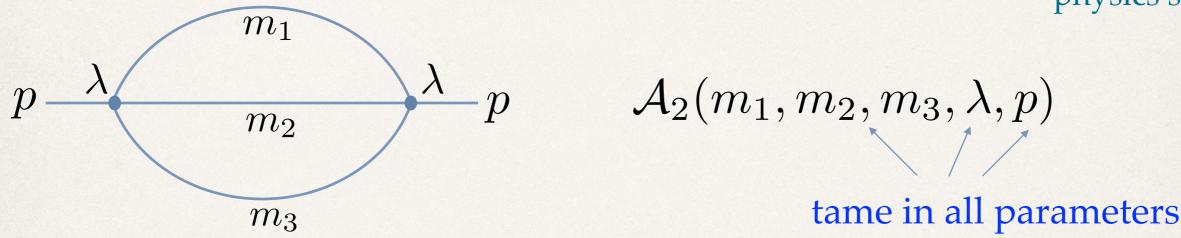


hidden finiteness property in all QFT amplitudes

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Remarks: - theorem is non-trivial: interesting implications for Feynman amplitudes (symmetry ← relations) [in progress]

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Tameness of (relative) periods

Tameness of Feynman integrals

Tameness in non-perturbative Quantum Field Theories

interested in physical observables in local QFTs

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 - particle described by the field ϕ
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(up to normalization)

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path integral over all field configurations exponential weight by parameter-dep. Lagrangian

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$$\langle \mathcal{O}_1(y_1)...\mathcal{O}_k(y_k)\rangle_{\lambda}$$

ightharpoonup complicated function on product of space-time $\Sigma \times ... \times \Sigma$ and parameter space $\mathcal P$

Consider in 0d:
$$S = \frac{m^2}{2}\phi^2 + \frac{\lambda}{4!}\phi^4 \rightarrow Z = \sqrt{\frac{3}{\lambda}}e^{\frac{3m^4}{4\lambda}} m K_{\frac{1}{4}}\left(\frac{3m^4}{4\lambda}\right)$$

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- → well-defined notion of complexity for physical systems

[TG,Schlechter, van Vliet] to appear

QFTs on finite lattice

correlation functions in 0d are ordinary integrals

$$\langle \mathcal{O}_1...\mathcal{O}_n \rangle_{\lambda} = \int d\phi_1...d\phi_k \,\mathcal{O}_1...\mathcal{O}_n \,e^{-S^{(0)}(\phi,\lambda)}$$

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Conjecture [van den Dries][Kaiser]: Given a real-valued tame function $f(\lambda,\phi)$ (in some o-minimal structure $\mathcal S$) the integral

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Note: Theorem for $S = \mathbb{R}_{an} \to \widetilde{S} = \mathbb{R}_{an,exp}$. [Comte,Lion,Rolin]

However, for non-perturbative results, we need exponential to be in \mathcal{S} .

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⇒ math. conjecture implies:

[Douglas, TG, Schlechter '23]

in order that physical observables $\langle \mathcal{O}_1...\mathcal{O}_n \rangle_{\lambda}$ are tame functions of parameters λ one needs to require:

 $S^{(0)}(\phi,\lambda)$ is tame function of λ,ϕ

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special, but interesting case: exponential periods with parameters

$$\Pi(\lambda) = \int_{\Gamma} e^{-f(\lambda)} \omega(\lambda) \qquad \qquad \frac{f(\lambda) \text{ algebraic function}}{\omega(\lambda) \text{ algebraic differential form}}$$

Question 3: Are exponential periods definable in $\mathcal{P}(\mathbb{R}_{an,exp})$?

Non-tameness of Lagrangian: easy to get non-tame Lagrangian by picking non-tame potential V(x)

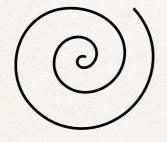
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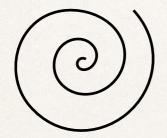
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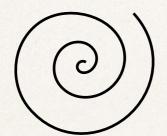
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More Fancy:
$$W_{\xi} = Y P_{\xi}(X_1, \dots, X_k)^2 + \sum_a Z_a (\sin 2\pi i X_a)^2$$
 [Tachikawa]

Existence of supersymmetric vacua is undecidable!

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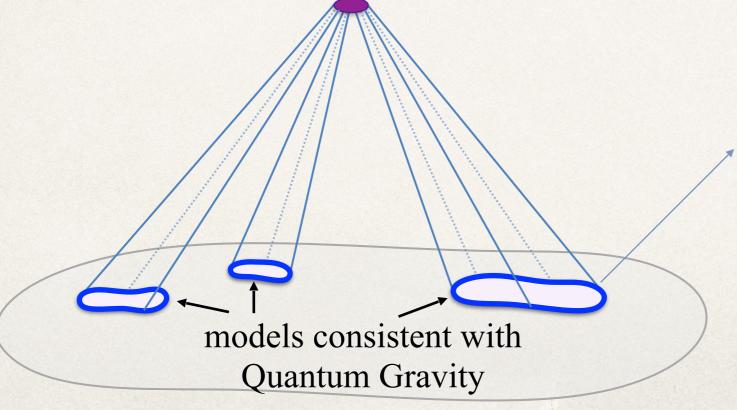
in general: tameness depends on the UV origin of the theory

Mapping out the tame parts of physics

- Tameness of Conformal Field Theory: precise conjectures [Douglas, TG, Schlechter '23]
 - (1) Correlation / partition functions are tame functions over Euclidean space-time and over parameter space.
 - (2) Space of CFTs is tame set under certain conditions (e.g. bound on degrees of freedom).

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- Tameness of effective theories compatible with Quantum Gravity



tame set of effective theories that has tame physical observables Thanks!

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Solutions: introduce cut-off $m \leq \Lambda_{\rm UV}$